Methane to Methanol Conversion Induced by Thorium Oxide through the $CH₃Th(O)H$ Intermediate in Solid Argon

Yu Gong,[†] Lester Andrews,^{*,†} Virgil E. Jackson,[‡] and David A. Dixon^{*,‡}

† Department of Chemistry, Unive[rsit](#page-4-0)y of Virginia, Charlottesville, Virginia 22904-4319, [Un](#page-4-0)ited States ‡ Department of Chemistry, The University of Alabama, Tuscaloosa, Alabama 35487-0336, United States

S Supporting Information

[AB](#page-4-0)STRACT: [Reactions of T](#page-4-0)hO molecules and CH4 have been investigated in solid argon near 4 K. The $CH₃Th(O)H$ molecule is produced when the sample is exposed to UV irradiation. Identification of this new intermediate is substantiated by observation of the Th=O and Th−H stretching vibrational modes with isotopic substitution via matrix infrared spectroscopy, and the assignments are supported by electronic structure frequency calculations. Methanol absorptions increase together with formation of the $CH₃Th(O)H$ molecule, suggesting a methane to methanol conversion induced by thorium oxide

proceeding through the CH₃Th(O)H intermediate. The formation of CH₃Th(O)H from ThO + CH₄ is exothermic (ΔH_{run} = −11 kcal/mol) with an energy barrier of 30 kcal/mol at the CCSD(T)//B3LYP level. Decomposition of this intermediate to form methanol involves spin crossing, and the overall reaction from the intermediate is endothermic by 127 kcal/mol. There is no activation energy for the reaction of thorium atoms with methanol to give $CH₃Th(O)H$, as observed in separate experiments with Th and CH₃OH.

NO INTRODUCTION

The conversion from methane to methanol has attracted numerous studies since methanol can be widely used as a fuel as well as the starting material for producing other useful chemicals.¹ The successful partial oxidation of methane requires effective catalysts to break the inert C−H bond. A large number of studie[s](#page-4-0) have been performed on the preparation and optimization of homogeneous, heterogeneous and enzyme catalysts as well as probing the catalytic mechanisms.^{2−5} In addition to these studies in the condensed phase, systematic gas phase studies using mass spectrometric methods have als[o](#page-4-0) [be](#page-4-0)en carried out in a wide range on the reactions of size-selected transition metal oxide ions toward methane as well as other hydrocarbons, as summarized in a series of recent reviews. $6-8$ Although these studies were focused on simplified model molecular systems, valuable information on the mechani[st](#page-4-0)i[c](#page-4-0) aspects can be obtained especially when combined with high level theoretical calculations.⁹ Apart from the numerous studies in the gas phase employing charged species, the reactions of neutral transition metal ox[id](#page-4-0)e molecules and methane have received less attention, and most results are from spectroscopic studies in noble gas matrixes.¹⁰ In particular, the cryogenic matrix environment makes it possible to relax and trap intermediates in the methane [o](#page-4-0)xidation reaction, which are usually not accessible in the gas phase.¹⁰

Although reactions of neutral transition metal oxide molecules and methane have been [s](#page-4-0)ystematically studied under matrix isolation conditions,^{10−12} no investigation has been carried out on the reactivity of neutral actinide oxide species, even though the role of a[ctinide](#page-4-0) elements in catalytic reactions has been the subject of many recent studies.^{13,14} The most easily accessible actinide metals, uranium and thorium, have been shown to activate the C−H bonds of [di](#page-4-0)[ff](#page-4-0)erent hydrocarbons.^{14,15} Gas phase mass spectrometric studies have provided detailed information on the reactivity and thermodynamics of act[inide](#page-4-0) containing cations.¹⁶ Recent experimental and theoretical studies on the reactions of thorium and uranium cations revealed that while thorium [c](#page-4-0)ations and dications reacted with CH₄ to form the ThCH₂⁺ and ThCH₂²⁺ products, no similar dehydrogenation reaction was observed for uranium cation owing to the endothermic character as well as the high energy barrier for the uranium reactions.^{17−21}

Our recent matrix isolation infrared spectroscopic studies have provided evidence for the activation [of](#page-4-0) [me](#page-4-0)thane by neutral uranium and thorium atoms, as the methylidene complexes $CH₂MH₂$ (M = U, Th) were observed. The computed structure of the uranium product is more distorted due to stronger agostic interactions. $22,23$ For reactions involving uranium and thorium oxide cations in the gas phase, the oxide−arene adducts were the [major](#page-4-0) products in most cases with H_2 and CH4 elimination observed in a few cases depending on the reactant arenes.16,24 The overall reaction efficiency for the actinide oxide cations was much lower than that for bare metal cations as evid[ence](#page-4-0)d by a series of reactivity studies with different hydrocarbons, and almost no reaction was observed between actinide oxide cations and methane.¹⁸ In the current work, we report the reactions of laser-ablated neutral thorium oxide molecules and methane in solid argon a[t 4](#page-4-0) K. Here, ThO molecules react with CH4 when broad band UV irradiation is employed, which produces the $CH₃Th(O)H$ intermediate trapped in a solid argon matrix. At the same time, evidence is also provided for the formation of methanol, suggesting that the

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oxidation of methane proceeds in the presence of thorium oxide molecules.

EXPERIMENTAL AND THEORETICAL METHODS

The experimental apparatus and procedure for studying laser-ablated ThO and CH₄ reactions has been described previously.^{25,26} The Nd:YAG laser fundamental (1064 nm, 10 Hz repetition rate with 10 ns pulse width) was focused onto a thorium dioxide (thoria [cera](#page-4-0)mic) target mounted on a rotating rod. Laser-ablated thorium oxide materials, (ThO, Th, and O) were co-deposited with 4 mmol of argon (Matheson, research grade) containing 0.5% CH₄ (Matheson) onto a CsI cryogenic window at 4 K. Isotopic precursors CD_4 and ¹³CH₄ (Cambridge Isotope Laboratories, 99% enrichment) were used without purification. Complementary Th and methanol experiments also employed $\text{CH}_3^{-18}\text{OH}$ (Cambridge Isotopic Laboratories, 95% enrichment). FTIR spectra were recorded at 0.5 cm⁻¹ resolution on a Nicolet 750 FTIR instrument with HgCdTe range B detector. Matrix samples were annealed at different temperatures and cooled back to 4 K for spectral acquisition. Selected samples were subjected to broadband photolysis by a medium-pressure mercury arc street lamp (Philips, 175 W) with the outer pyrex globe removed to emit the entire quartz envelope output (220 < λ < 900 nm), which was also reduced using selected long wavelength pass optical glass filters.

Density functional theory $(DFT)^{27}$ with the B3LYP hybrid exchange–correlation functional²⁸ was used to optimize all of the structures. The initial calculations on ${}^{1}CH_{3}Th(O)H$ ${}^{1}CH_{3}Th(O)H$ ${}^{1}CH_{3}Th(O)H$ were done with the 6-311++G(d,p) basis set²⁹ o[n H](#page-4-0)₁, C, and O and the SDD effective core potential and basis set on Th. 30 (The preceding superscript in a formula corresponds to the [u](#page-5-0)sual spin state definition.) Subsequent calculations at the DFT/B3LYP lev[el](#page-5-0) with the aug-cc-pVTZ 31 basis set on H, C, and O and the small core relativistic effective core potential (ECP) from the Stuttgart group30,32 with the corr[es](#page-5-0)ponding segmented [8s,7p,6d,4f,2g] basis set on Th were used to map out the potential energy surface. The G[aussi](#page-5-0)an 09 program system was used for the optimization and frequency calculations of the minima and for the final DFT energies. 33 We were able to optimize some of the transition states using the options available for such optimizations in the Gaussian code.³⁴ In cas[es](#page-5-0) where we could not determine the transition state, we used the STEPPER approach with specific mode following in NWChem.³⁵ The DFT geometries were used in single point CCSD(T) calculations³⁶ performed with the MOLPRO 2010.1 program system³⁷ and [wi](#page-5-0)th open shell molecules treated at the R/ $UCCSD(T)$ level.³⁸ The sa[me](#page-5-0) ECP for Th was used as given above, but a larger basi[s s](#page-5-0)et of the form [10s,9p,5d,4f,3g] was used following our previous work [on](#page-5-0) ThO_2^{39} with the aug-cc-pVTZ basis set on H, C, and O. We correlated³⁹ 12 valence electrons on Th in the $CCSD(T)$ calculations (nominally the [6s](#page-5-0), 6p, 7s, and 6d) and 6 electrons on O, 4 electrons on C, and [1 e](#page-5-0)lectron on each H atom.

■ RESULTS AND DISCUSSION

Infrared spectra from the reactions of laser-ablated thorium oxide materials (ThO, Th, and O) and methane in solid argon are shown in Figure 1. Strong absorption due to ThO, as well as weak $ThO₂$ bands were observed right after sample deposition (Figure 1, trace a), which result from laser ablation of the thorium oxide target. The yield of ThO in these experiments is much higher than that of $ThO₂$ relative to previous work with Th + O_2 reactions,⁴⁰ so we must conclude that either some ThO molecules survive ablation or that extensive recombination of the element[s g](#page-5-0)ives the high yield of ThO. In addition, the broad methyl radical absorption centered at 603 cm[−]¹ (not shown here) was observed, 41 which is common in all of our reactions of laser-ablated metal atoms and methane, where the laser ablation plume ph[oto](#page-5-0)dissociates methane.²⁶ Trace absorptions were observed for two matrix trapping sites of CH₃OH at 1033.7 and 1027.0 cm⁻¹ on sample d[ep](#page-4-0)osition: these likely arise from reaction of laser-ablated O atoms with

Figure 1. Infrared spectra of laser-ablated ThO, Th, and O and CH₄ reaction products in solid argon: (a) ThO + 0.5% CH₄ deposition for 60 min; (b) after annealing to 20 K; (c) after λ >290 nm irradiation; (d) after $\lambda > 220$ nm irradiation; (e) annealing to 30 K; (f) after annealing to 35 K.

methane. Also weak bands at 1435.7, 1397.1, and 670.8 cm^{-1} due to the $CH₂ThH₂$ methylidene complex result from the reactions of thorium atoms and methane.²² Subsequent sample annealing to 20 K increased the $ThO₂$ absorptions as well as a band at 842.6 cm[−]¹ (Figure 1, trace b)[, w](#page-4-0)hich was assigned previously to the HThO molecule.⁴² The CH_2THH_2 bands increased and a new band pair was produced at 1370 and 817 cm⁻¹ on exposure to pyrex filtered ($\lambda > 290$ nm) irradiation. When the sample was next exposed to full arc $(\lambda > 220 \text{ nm})$ irradiation, the CH_2ThH_2 bands and the new pair at 1370 and 817 cm⁻¹ increased along with the 1027.0 cm⁻¹ matrix site band of the C−O stretching mode for CH₃OH. Together with the absorptions due to CH_2ThH_2 , the 1370 and 817 cm⁻¹ absorptions sharpened and shifted to 1375.5 and 816.7 cm^{-1} during sample annealing to 30 K but decreased slightly upon sample annealing to 35 K. After the absorption due to methanol increased at 1027.0 cm[−]¹ upon UV irradiation, conversions between the two 1027.0 and 1033.7 cm[−]¹ matrix site absorptions occurred during sample annealing (see Figure 1, traces d and e). Clearly, the increase of CH₃OH on full arc (λ > 220 nm) irradiation accompanies a further increase of the new pair of bands to be identified here.

Complementary studies of the reaction of thorium and methanol were also performed, and these showed similarities with the thorium oxide and methane reaction (Figure 2, trace d). The new band at 816.7 cm[−]¹ appeared as well when the sample was annealed after sample deposition. Howe[ve](#page-2-0)r, the 1375.5 cm[−]¹ band in the thorium and methanol experiment was not well resolved because of overlap with the methanol precursor absorptions.

A new product molecule resulting from the reaction of thorium oxide and methane must be considered since the 1375.5 and 816.7 cm[−]¹ absorptions were not observed in reactions of thorium and methane. The 816.7 cm[−]¹ band should be due to a terminal Th=O stretching mode due to the band position as well as the characteristic $16O/18O$ ratio of 1.0560 obtained from the Th and $\text{CH}_3^{18}\text{OH}$ reaction (Figure 2). Experiments with ¹³CH₄ revealed no ¹³C shift for the 816.7 cm^{-1} band, and the reaction of thorium oxide and CD_4 resulted [in](#page-2-0) a very small red shift of 1.8 cm[−]¹ , which is comparable with

Figure 2. Infrared spectra of laser-ablated ThO, Th, and O and isotopically substituted CH_4 reaction products in solid argon (all spectra taken after $\lambda > 220$ nm irradiation followed by annealing to 30 K): (a) ThO + 0.5% CH₄; (b) ThO + 0.5% ¹³CH₄; (c) ThO + 0.5% CD4. (d, e) Infrared spectra of laser-ablated Th atoms and isotopically substituted CH₃OH reaction products in solid argon: (d) Th + 1.0% CH₃OH, after annealing to 35 K; (e) Th + 0.5% CH₃¹⁸OH, after annealing to 40 K.

the 2.1 cm⁻¹ red shift for the HThO molecule (Figure 2).⁴² These observations suggest that this new terminal $Th=O$ stretching mode is slightly perturbed by vibrational motio[ns](#page-5-0) involving hydrogen. For the other band at 1375.5 cm[−]¹ , an experiment with $^{13}CH_4$ gave no isotopic shift, indicating that carbon in the new molecule is not involved in this higher frequency vibrational mode. When $CD₄$ was used as the reactant, no band was observed between 1370 and 1450 cm⁻¹. . This follows since absorptions in this region have been assigned to Th–H stretching modes of the CH_2THH_2 molecule.²² Hence it is also reasonable to assign the 1375.5 cm⁻¹ band to the Th−H stretching mode of the new molecule. If the H[/D](#page-4-0) isotopic ratio for the 1375.5 cm[−]¹ band is similar to that of other Th–H containing species $(1.39-1.40)$,⁴² the deuterium counterpart would be observed in the region of 980−990 cm[−]¹ ; this is on the red side of the strong CD_4 [abs](#page-5-0)orption, which broadens upon sample annealing. On the basis of the two experimental frequencies, the new product molecule is identified as $CH_3Th(O)H$ with terminal Th= O and Th−H moieties.

Our assignment of this new molecule is strongly supported by calculations at the DFT/B3LYP level (Table 1). The $CH₃Th(O)H$ molecule is predicted to have a closed shell singlet ground state with a pyramidal geometry (Figure 3), similar to the recently characterized $CH₃U(O)H$ molecule.⁴³ The $Th=O$ bond length and the corresponding stretching frequency for $CH_3Th(O)H$ are also close to those of the t[wo](#page-5-0) previously characterized H_2ThO and $(CH_3)_2ThO$ molecules with $Th(IV)$ centers.^{42,44} The calculations predict two bands at 1397.2 and 834.8 cm[−]¹ that are much stronger than all of the

Figure 3. Structure of the ^{1}A ground state of the CH₃Th(O)H molecule calculated at the B3LYP level of theory. Bond lengths in Å. Top numbers using the $6-311++G(d,p) + SDD$ basis set. Bond lengths and angles with the aug-cc-pVTZ + ECP basis set in parentheses: ∠OThH, 104.6° (104.8°); ∠CThH, 96.8° (96.8°); ∠CThO, 103.1° (103.2°); ∠sum = 304.5° (304.8°). Th is blue, O is red, C is gray, and H atoms are white.

other absorptions above 400 cm[−]¹ . The Th−H harmonic stretching band predicted at 1397.2 cm^{-1} is slightly higher than the observed anharmonic frequency at 1375.5 cm[−]¹ . The deuterium counterpart is calculated to be 993.4 cm[−]¹ with only half of the intensity of the 1397.2 cm⁻¹ band, in the region where the CD₄ precursor molecule absorbs. The Th-D band for the $CD_3Th(O)D$ molecule is probably covered by the strong CD4 precursor band as noted above. For the calculated Th= \overline{O} stretching mode at 834.8 cm⁻¹, the ¹⁶O/¹⁸O ratio is 1.0562. Both the band position and the isotopic frequency ratio agree well with the experimental results. We note that the calculated values are for the harmonic frequencies and the experimental values include anharmonic components as well as the influence of matrix atoms. These values reported to 0.1 cm[−]¹ come from precise calculations and experiments, but one should not associate an error bar with the calculated frequencies to this accuracy.

In order to explore the reaction mechanism, the potential energy surface for the ¹ThO + CH₄ \rightarrow ³Th + CH₃OH reaction was calculated at the CCSD(T)//B3LYP level of theory (Figure 4) with the CCSD(T) energies including the B3LYP zero point energy correction. All of the singlet and triplet species [ex](#page-3-0)cept for the $\mathrm{^{1}ThO}(\mathrm{CH}_4)$ complex discussed below are shown. Formation of the $^{1}CH_{3}Th(O)H$ molecule from the reaction of ${}^{1}\text{ThO}$ and CH_4 is predicted to be exothermic $(\Delta H_{\text{rxn}} = -11 \text{ kcal/mol})$, with an energy barrier (¹TS1) of 30 kcal/mol separating reactants and the product. This modest energy barrier is consistent with the appearance of the ¹CH₃Th(O)H absorptions under broadband mercury arc UV irradiation. Geometry optimization of the $\mathrm{^{1} ThO(CH_{4})}$ complex resulted in separated ${}^{1}\text{ThO}$ and CH_4 at the DFT and MP2 levels, in agreement with the absence of the observation of this molecular complex in our experiments. ${}^{3}CH_{3}Th(O)H$ is predicted to be 56 kcal/mol above $^{1}CH_{3}Th(O)H$. Further decomposition of $^{1}CH_{3}Th(O)H$ is prohibited by energy barriers of 47 and 77 kcal/mol separating it from $^{1}CH_{3}ThOH$ and ¹CH₃OThH with the former being 24 kcal/mol more

Table 1. Observed and Calculated (in Bold) Infrared Absorptions $\rm (cm^{-1})$ of the $\rm CH_3Th(O)H$ Molecule in Solid Argon a

 a Numbers in parentheses are infrared intensities in km/mol calculated at B3LYP level of theory. b Not observed due to band overlap. c When the larger aug-cc-pVTZ(ECP) basis set at the B3LYP level is used, the Th−H frequency is 1394.4 (477) cm⁻¹ and the Th=O frequency is 845.2 (255) cm^{-1} . .

Figure 4. Calculated potential energy surface for the $^1\text{ThO} + \text{CH}_4 \to {}^3\text{Th} + \text{CH}_3\text{OH}$ reaction at the $\text{CCSD(T)}/\text{aug-cc-pVTZ(ECP)}/\text{B3LYP}/\text{aug-cc-}\text{poly}$ cc-pVTZ(ECP') (see text for basis set details) level of theory in kcal/mol. The imaginary frequencies for the transition states are 1315i ($\rm ^1TS1)$, 688i (³TS1), 670i (¹TS2′), 1091i (³TS2′), 1439i (¹TS2″), 466i (³TS2″), 498i (³TS3′), and 834i (³TS3″). The spin–orbit corrected value for the ³Th + ¹CH OH spin–orbit corrected value for the ³Th + CH₃OH asymptote is in parentheses, and that for the ${}^{3}\text{Th}\cdots\text{CH}_{3}\text{OH}$ complex is 81 kcal/mol.

stable than the latter. Both $^{1}CH_{3}ThOH$ and $^{1}CH_{3}OThH$ are predicted to be less stable than the reactant asymptote of ThO + CH₄. As observed in other divalent thorium species, $42,44,45$ both the $CH₃OThH$ and $CH₃ThOH$ molecules have singlet ground states. ${}^{3}CH_{3}ThOH$ is predicted to be 14 kc[al/mol](#page-5-0) higher than 1 CH₃ThOH and 3 CH₃OThH is predicted to be 12 kcal/mol above $^{1}CH_{3}OH$ H. The presence of relatively small singlet−triplet splittings for Th (a heavy element) should enable spin orbit coupling in this region and crossing from the singlet to the triplet thereby allowing access to the lowest energy 3 Th + 1 CH₃OH product asymptote. The molecular singlet–triplet splitting is smaller than the ³F-¹D splitting of 20.8 kcal/mol in the thorium atom, which is consistent with this hypothesis.⁴⁶

The $3\text{Th} + 1\text{CH}_3\text{OH}$ asymptote is predicted to be 128 kcal/ mol above the [rea](#page-5-0)ctant asymptote ${}^{1}ThO + CH_{4}$. Application of a spin orbit correction to the ³Th atom ($\Delta E_{\textrm{SO}} = \sum_J [(2J +$ $1|\Delta E(J)|\sum_{I}[(2J + 1)]$) using the experimental spin orbit splittings⁴⁶ for the Th atom $(\Delta E(J))$ gives a reaction energy of 119 kcal/mol, which is in excellent agreement with the value of 118.7 \pm [3.8](#page-5-0) kcal/mol calculated from the experimental heats of formation.⁴⁷ A stable ³Th…CH₃OH complex is 38 kcal/mol lower in energy than the product asymptote. This complex is separated [by](#page-5-0) a 22 kcal/mol barrier from ${}^{3}CH_{3}OH_{1}$, which is the higher energy isomer. The moderate barrier is due to the ease of hydrogen transfer. A much higher barrier of 58 kcal/mol separates the 3 Th…CH₃OH complex from the more stable isomer ${}^{3}CH_{3}ThOH$. The higher barrier is due to the difficulty in transferring the CH₃ group and actually lies above the 3 Th + 1 CH OH product asymptote. Since the singular reacture ${}^{1}CH_{3}OH$ product asymptote. Since thorium atoms react spontaneously with methanol, it is most likely for the reaction to proceed through the higher energy ${}^{3}CH_{3}OH$ isomer, which essentially requires no activation energy. A similar mechanism has been observed in the reactions of scandium and

lanthanide atoms with methanol.⁴⁸ Our proposed pathway through the higher energy isomer with the lower barrier is also consistent with the fact that $CH₃OH$ can be formed by broadband UV irradiation (Figure 1, trace d), where the UV energies applied in the experiment have approximately zero intensity at 220 nm (130 kcal/mol[\) a](#page-1-0)nd considerable intensity at 240 nm (119 kcal/mol). The energy of the 240 nm photon is just enough to overcome the endothermicity of the reaction from the stable intermediate to ${}^{3}Th + CH_3OH$, and both photon energies are too low for the highest energy barrier to be accessed. This is consistent with the fact that lower energy $>$ 290 nm photolysis produces the CH₃Th(O)H intermediate but not additional CH₃OH. All of the molecules on the potential energy surface prefer singlet ground states except for the ³Th + CH₃OH asymptote and the ³Th(CH₃OH) complex. This suggests that intersystem crossing occurs during the overall reaction, which is consistent with gas phase reactions involving transition metal oxide cations with high spin ground states.^{6d,49}

Previous studies on the reactions of transition metal oxide mole[cul](#page-4-0)[es](#page-5-0) and methane revealed that $CH₃M(O)H$ intermediates were observed when laser-ablated TiO, NbO, and TaO molecules reacted with methane whereas the reactions of MnO and FeO with CH_4 gave the CH₃MOH products without metal−oxygen multiple bonds.¹⁰ Formation of the CH₃Th(O) H molecule further supports the similarity of thorium with the group IV metals, both of whic[h h](#page-4-0)ave a highest oxidation state of IV, as is also displayed by the corresponding dioxides.^{40,50} However, no $ThO(CH_4)$ complex was observed in the thorium oxide experiments, which differs from the titanium cas[e. It](#page-5-0) should be noted that TiO possesses a triplet ground state whereas the closed shell singlet is the most stable for the ThO molecule.⁵¹ There is almost no reaction between ThO⁺ cations and CH_4 from previous gas phase studies.¹⁸ This reaction

profile should be different from that of the neutrals because the $CH₃Th(O)H⁺$ cation will have a weaker Th–O bond energy with only three electrons available for that interaction rather than the four electrons available for neutral $CH₃Th(O)H$. Hence, the CH₃ThOH⁺ cation with Th–C and Th–OH bonds or its decomposition products are expected to be formed, as observed in gas phase studies on the reactions of transition metal oxide cations with methane.⁶

■ CONCLUSION

Matrix isolated ThO molecules react with $CH₄$ to give the CH3Th(O)H product in solid argon at 4 K during UV irradiation. Infrared absorptions due to methanol also increased, along with the formation of $CH₃Th(O)H$ using full arc $(\lambda > 220 \text{ nm})$ irradiation, suggesting a methane to methanol conversion induced by thorium oxide going through the $CH₃Th(O)H$ intermediate. Calculations at the B3LYP level reveal that the $CH₃Th(O)H$ molecule possesses a pyramidal structure with a closed shell singlet ground state. From the calculated potential energy surface, formation of the $CH₃Th-$ (O)H molecule from the reaction of ThO and methane has an energy barrier of 30 kcal/mol and the $CH₃Th(O)H$ product is 11 kcal/mol more stable than the asymptotic reactants. Decomposition of this intermediate to $3\text{Th} + \text{CH}_3\text{OH}$ is endothermic by 127 kcal/mol, and the formation of ³Th will require a spin crossing. On the other hand, no activation energy is required for the reaction of ³Th and methanol, which results in the spontaneous production of ${}^{1}CH_{3}Th(O)H$, again with a surface crossing needed. The reactions of thorium oxide and methane studied here suggest that actinide oxides such as ThO might be used to oxidize methane to methanol under photochemical conditions, analogous to transition metal oxide molecules.10,11

■ ASSOCIATED CONTENT

6 Supporting Information

Complete references 33 and 37 and tables containing thermochemical values at the CCSD(T) and B3LYP level of theory, frequencies and [IR](#page-5-0) inten[sitie](#page-5-0)s at the B3LYP level of theory for ${}^{1}CH_{3}Th(O)H$, ${}^{1}CH_{3}ThOH$, and ${}^{1}CH_{3}OH$ and optimized B3LYP Cartesian coordinates for the various structures. This material is available free of charge via the Internet at http://pubs.acs.org.

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Corresponding Author

*E-mail: lsa@virginia.edu (L.A.); dadixon@as.ua.edu (D.A.D.).

Notes

The auth[ors declare no c](mailto:lsa@virginia.edu)ompeting fi[nancial interest.](mailto:dadixon@as.ua.edu)

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